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Einstein metrics with A_{16} -type holonomy

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Abstract

Einstein neutral metrics in dimension four are constructed. They provide examples of holonomy type A_{16} .

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The holonomy group of a metric g at a point x of a manifold M is the group of linear transformations in the tangent space of x defined by parallel translation along all possible loops starting at x . For connections on connected manifolds, holonomy groups of different points are isomorphic, and so we shall refer to *the* holonomy group of g and denote it by $Hol(g)$ [1].

If one restricts to curves which are null homotopic, one obtains the restricted holonomy group $Hol^\circ(g)$ of M , which is the identity component of $Hol(g)$. Clearly $Hol(g)$ and $Hol^\circ(g)$ are equal if M is simply connected. Also, $Hol(g)$ and $Hol^\circ(g)$ are Lie groups.

It is obvious that a connection can only be the Levi-Civita connection of a metric g if the holonomy group is a subgroup of the generalized orthogonal group corresponding to the signature of g [2].

At any point $x \in M$, and in some coordinate system about x , the set of matrices of the form

$$\mathbf{R}_{bcd}^a X^c Y^d \quad \mathbf{R}_{bcd:e}^a X^c Y^d Z^e \quad \mathbf{R}_{bcd:ef}^a X^c Y^d Z^e W^f \dots$$

where $X, Y, Z, W \in T_x M$ and a semi-colon denotes a covariant derivative, forms a Lie subalgebra of the Lie algebra of $M_n(\mathbb{R})$ of $GL(n, \mathbb{R})$ called the *infinitesimal holonomy algebra* of M at x . Up to isomorphism the latter is independent of the coordinate system chosen and is denoted by $hol'(g)$. The corresponding uniquely determined connected subgroup of $GL(n, \mathbb{R})$ is called the infinitesimal holonomy group of M at x and is denoted by $Hol'(g)$. The Lie algebra $hol'(g)$ is a subalgebra of the Lie algebra of $Hol(g)$ for each $x \in M$ [3–5].

In a recent paper, Ghanam and Thompson [6] studied and classified the holonomy Lie subalgebras of neutral metrics in dimension four. In this paper, we will study one of these subalgebras A_{16} and construct Einstein metrics of holonomy type A_{16} .

A_{16} is a two-dimensional Lie subalgebra generated by $\begin{bmatrix} -J & 0 \\ 0 & J \end{bmatrix}$ and $\begin{bmatrix} J & 0 \\ 0 & J \end{bmatrix}$ where $J = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$.

By a change of basis, one obtains a new set of generators, namely, $\begin{bmatrix} J & 0 \\ 0 & 0 \end{bmatrix}$ and $\begin{bmatrix} 0 & 0 \\ 0 & J \end{bmatrix}$. Consider \mathbb{R}^2 with coordinates (x, y) , and let

$$g = \begin{bmatrix} a & c \\ c & b \end{bmatrix} \quad (1)$$

be a two-dimensional metric, where a, b and c are smooth functions of x and y . Then the equation $gJ + (gJ)^t = 0$ implies that $a = b$ and $c = 0$. Hence $g = a(x, y)I$, where I is the 2×2 identity matrix. We apply the de Rham decomposition theorem: use a pair of non-flat two-dimensional Riemannian metrics with the signs adjusted so as to produce a neutral four-dimensional metric. This proves that A_{16} is a holonomy Lie algebra of four-dimensional neutral metric.

Let us turn to equation (1). The non-zero components of the Ricci tensor are

$$R_{11} = R_{22} = \frac{(a_{yy}a + a_{xx}a - a_y^2 - a_x^2)}{2a^2} \quad (2)$$

and the Ricci scalar is given by

$$R = \frac{(a_{yy}a + a_{xx}a - a_y^2 - a_x^2)}{a^3}. \quad (3)$$

The Einstein condition entails that

$$R_{ij} = \frac{1}{4}Rg_{ij} \quad (4)$$

and so if we take $i = j = 1$ in equation (4), we obtain the following PDE:

$$a(a_{xx} + a_{yy}) - (a_x^2 + a_y^2) = 0. \quad (5)$$

We will solve equation (5) by separation of variables. Let

$$a(x, y) = f(x)h(y) \quad (6)$$

where f and h are smooth functions of x and y , respectively. Equation (5) becomes

$$fh(f''h + fh'') - ((f')^2h^2 + f^2(h')^2) = 0 \quad (7)$$

and so

$$h^2(ff'' - (f')^2) + f^2(hh'' - (h')^2) = 0. \quad (8)$$

If we divide equation (8) by f^2h^2 we obtain

$$\frac{ff'' - (f')^2}{f^2} = \frac{(h')^2 - hh''}{h^2}. \quad (9)$$

This implies that

$$\frac{ff'' - (f')^2}{f^2} = \frac{(h')^2 - hh''}{h^2} = c \quad (10)$$

where c is a constant. Now we will consider two cases:

Case (1). If $c = 0$, then equation (10) gives

$$ff'' - (f')^2 = 0 \quad (11)$$

and

$$(h')^2 - hh'' = 0. \quad (12)$$

To solve equation (11) we use the following substitution:

$$f(x) = e^{z(x)} \quad (13)$$

where $z(x)$ is a smooth function of x . We substitute equation (13) in equation (11) to obtain

$$e^z(z''e^z + (z')^2e^z) - (z')^2e^{2z} = 0 \quad (14)$$

and so

$$z'' = 0. \quad (15)$$

Hence

$$z(x) = c_1x + c_2 \quad (16)$$

where c_1 and c_2 are constants. Hence

$$f(x) = e^{c_1x+c_2}. \quad (17)$$

Similarly equation (12) implies that

$$h(y) = e^{c_3y+c_4} \quad (18)$$

where c_3 and c_4 are constants. Therefore the two-dimensional Einstein metric g is given by

$$g = e^{c_1x+c_2y+c_3} (dx^2 + dy^2) \quad (19)$$

and so the four-dimensional Einstein neutral metric with holonomy type A_{16} is

$$g = e^{c_1x+c_2y+c_3} (dx^2 + dy^2) - e^{c_4z+c_5t+c_6} (dz^2 + dt^2) \quad (20)$$

where (x, y, z, t) is a coordinate system on \mathbb{R}^4 .

Case (2). If $c \neq 0$, then equation (10) gives

$$ff'' - (f')^2 - cf^2 = 0 \quad (21)$$

and

$$(h')^2 - hh'' - ch^2 = 0. \quad (22)$$

To solve equation (21), we use equation (13) to obtain

$$z'' - c = 0. \quad (23)$$

and so

$$z(x) = \frac{1}{2}cx^2 + c_1x + c_2. \quad (24)$$

To solve equation (22), we let $h(y) = e^{w(y)}$ to obtain

$$w'' + c = 0 \quad (25)$$

and so

$$w(y) = -\frac{1}{2}cy^2 + c_3y + c_4. \quad (26)$$

Thus, the two-dimensional Einstein metric g is given by

$$g = e^{\frac{1}{2}c(x^2-y^2)+c_1x+c_2y+c_3} (dx^2 + dy^2) \quad (27)$$

and so the four-dimensional Einstein neutral metric with holonomy type A_{16} is

$$g = e^{\frac{1}{2}c(x^2-y^2)+c_1x+c_2y+c_3} (dx^2 + dy^2) - e^{\frac{1}{2}r(z^2-t^2)+c_4z+c_5t+c_6} (dz^2 + dt^2). \quad (28)$$

where $c_1, c_2, c_3, c_4, c_5, c_6$ and r are constants.

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